

Comments on the Mean Diffusion Length of Luminescent Molecules in Solutions *

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The effect of concentration of donor and acceptor molecules on the mean diffusion length in solutions is discussed.

The influence of mass diffusion on the non-radiative electronic excitation energy transfer (NEET) in luminescent solutions has been the subject of numerous publications [1–3].

In the discussion of this problem several zones have been distinguished limited by the critical distance R_{0A} for the NEET from D^* to A and by the mean diffusion length \bar{r}_0 , defined as [1]

$$\bar{r}_0 = (2\mathcal{D}\tau_{0D})^{1/2}. \quad (1)$$

In (1) τ_{0D} is the mean decay time of fluorescence (FL) of the donor D in the presence of the acceptor A , and \mathcal{D} equals $\mathcal{D}_D + \mathcal{D}_A$, where \mathcal{D}_D and \mathcal{D}_A are the diffusion coefficients of D and A , respectively.

It follows from theoretical considerations that mass diffusion should increase the efficiency of NEET [3, 4]. Results of experimental investigations are ambiguous. Some authors report a considerable influence of diffusion on NEET [5, 6], while only a slight or even negligible effect has been observed by others [7, 8].

Such discrepancies might have arisen for many reasons [3]. In our opinion, they may also be the result of considering the diffusion length \bar{r} to be a constant independent of the concentration of D and A in the solution.

$$\varphi(\gamma, \alpha, \alpha_0) = \frac{1 + \gamma^2 - (1.5 + \gamma^2)f + \alpha_0\alpha f(f + \gamma^2 - 0.5) - \alpha_0\alpha\gamma^2}{1 - f}; \quad \alpha = \gamma_D/\gamma. \quad (5)$$

α_0 is the probability of the excitation energy not being degraded during its transfer between donor molecules.

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Relation (4) is regarded as more general than formula (3). Indeed, if monomer quenching is excluded

Values of \bar{r}_0 as calculated according to (1) are correct only for $R_{DA} \gg R_{0A}$, where R_{DA} is the mean distance between D^* and A . If, however, this inequality is not satisfied at higher concentrations, then the mean diffusion length \bar{r} can be much shorter than \bar{r}_0 due to the additional deactivation channel for D^* molecules by NEET from D^* to D and A .

For sufficiently high concentrations the time of localization of the excitation energy on a D^* molecule,

$$\tau_1 = \left(\sum_D k_{D^*D} + \sum_A k_{D^*A} + k_F + k_q \right)^{-1},$$

can be much shorter than the decay time τ_{0D} (k_{D^*D} , k_{D^*A} , k_F and k_q are the rate constants for NEET from D^* to D , D^* to A , FL emission and internal conversion, respectively). In such a case the actual diffusion length should be calculated from relation (1) by substitution of τ_{0D} with τ_1 , i.e.

$$\bar{r} = (2\mathcal{D}\tau_1)^{1/2}. \quad (2)$$

The localization time τ_1 has been determined [9, 10] to be

$$\tau_1 = \tau_{0D} \frac{1 + \gamma^2 - (1.5 + \gamma^2)f}{1 - f}; \quad (3)$$

$$\gamma = \frac{1}{2} \pi^{1/2} (C_D/C_{0A} + C_A/C_{0A}) = \gamma_D + \gamma_A,$$

$$f \equiv f(\gamma) = \pi^{1/2} \gamma \exp(\gamma^2) [1 - \text{erf}(\gamma)].$$

Recently a more exact relation has been derived [11]:

$$\tau_1 = \tau_{0D} \varphi(\gamma, \alpha, \alpha_0) \quad (4)$$

where

($\alpha_0=1$) then for $\gamma_A \gg \gamma_D$, $\alpha \rightarrow 0$ and relation (4) becomes identical with (3). For $\gamma_D \gg \gamma_A$ relation (4) has a particularly simple form

$$\tau_1 = \tau_{0D}(1 - f). \quad (6)$$

Substitution of (4) into (2) leads to

$$\bar{r} = \bar{r}_0 [\varphi(\gamma, \alpha, \alpha_0)]^{1/2}. \quad (7)$$

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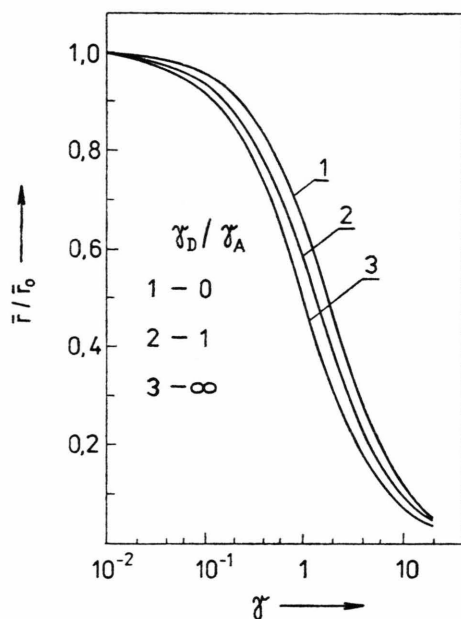


Figure 1 shows changes of \bar{r} as dependent on the reduced concentration γ for several values of α . Curves 1 and 3 correspond to the limiting values $\alpha = 0$ and $\alpha = 1$. Curve 2 was calculated for $\gamma_D = \gamma_A$, ($\alpha = 1/2$). Comparison of curves 1 and 3 shows that large variations of the ratio γ_D/γ_A have a relatively small effect on the concentration dependence of τ_1 , and thus \bar{r} . It is evident from the figure that the actual diffusion length \bar{r} can for high concentrations be much smaller than \bar{r}_0 . For $\gamma = 1$, ($C_A \approx C_{0A}$), \bar{r} is already equal to approximately $\bar{r}_0/2$, while for $\gamma = 10$, \bar{r} drops by more than one order of magnitude (see curve 1). It is thus possible that the discussion based on the value \bar{r}_0 instead of r can lead in some cases to the erroneous conclusion that mass diffusion does not affect the NEEET process in solutions.

Fig. 1. Actual diffusion length \bar{r} vs reduced concentration for different values of γ_D/γ_A and $\alpha_0 = 1$, calculated from (7).

- [1] J. B. Birks, *Photophysics of Aromatic Molecules*, John Wiley Interscience, New York 1970.
- [2] I. B. Berlman, *Energy Transfer Parameters of Aromatic Compounds*, Academic Press, New York 1973.
- [3] V. L. Ermolaev, I. N. Bodunov, J. B. Sveshnikova, and T. A. Shahverdov, *Nonradiative Electronic Excitation Energy Transfer* (in Russian), Science, Leningrad 1977.
- [4] I. Steinberg and E. Katchalski, *J. Chem. Phys.* **48**, 2404 (1968).
- [5] W. H. Melhuish, *J. Phys. Chem.* **67**, 1681 (1963).
- [6] Z. Varkonyi, *Acta Phys. et Chem., Szeged* **15**, 19 (1969).
- [7] Th. Förster, *Z. Elektrochem.* **53**, 93 (1949).
- [8] G. Porter and C. J. Tredwell, *Chem. Phys. Letters* **56**, 278 (1978).
- [9] C. Bojarski and J. Dudkiewicz, *Chem. Phys. Letters* **67**, 450 (1979).
- [10] C. Bojarski and E. Grabowska, *Acta Phys. Hung.* **40**, 113 (1979).
- [11] R. Twardowski and C. Bojarski, *Z. Naturforsch.* **35a**, 345 (1980).